High-Mass Star Formation

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The Importance of Massive Stars

McCaughrean+

Abel+

Vogelsberger+

Whitmore+

Gillessen+

Zinnecker & Yorke (2007) Tan et al. (2014) O'Dell+

The Physics of High-Mass Star Formation

A complicated, nonlinear process:

- Gravity vs pressure (thermal, magnetic, turbulence, radiation, cosmic rays) and shear.

 Heating and cooling, generation and decay of turbulence, generation (dynamo) and diffusion of B-fields.

- Chemical evolution of dust and gas.
- Fragmentation
- Stellar structure and evolution
- Feedback

- Wide range of scales (~12 dex in space, time) and multidimensional.

- Uncertain/unconstrained initial conditions/boundary conditions.



Complete theory of star formation

(Massive) Star Formation: Open Questions

- Causation: external triggering or spontaneous gravitational instability?
- Initial conditions: how close to equilibrium?
- Accretion mechanism: [turbulent/magnetic/thermalpressure]-regulated fragmentation to form cores vs competitive accretion / mergers
- Timescale: fast or slow (# of dynamical times)?
- End result
 - -Initial mass function (IMF)
 - -Binary fraction and properties



How do these properties vary with environment? Subgrid model of SF? Threshold n_{H*}? Efficiency ε_{ff}?

Outline

- Environments of Massive Star Formation
- Initial Conditions
- Timescales and Infall Rates
- Protostars Accretion & Outflow
- Feedback
- Dynamical Interactions











Massive Star Formation Theories

Core Accretion:

wide range of dm*/dt ~10⁻⁵ - 10⁻² M_{\odot} yr⁻¹

(e.g. Myers & Fuller 1992; Caselli & Myers 1995; McLaughlin & Pudritz 1997; Osorio+ 1999; Nakano+ 2000; Behrend & Maeder 2001)

Turbulent Core Model:

(McKee & Tan 2002, 2003) Stars form from **"cores"** that fragment from the **"clump"**



 $\bar{P} = \phi_P G \Sigma^2$

If in equilibrium, then self-gravity is balanced by internal pressure: B-field, turbulence, radiation pressure (thermal P is small)

Cores form from this turbulent/magnetized medium: at any instant there is a small mass fraction in cores. These cores collapse quickly to feed a central disk to form individual stars or binaries.

 $\dot{m}_* \sim M_{\rm core}/t_{\rm ff}$

Competitive (Clump-fed) Accretion:

(Bonnell, Clarke, Bate, Pringle 2001; Bonnell, Vine, & Bate 2004; Schmeja & Klessen 2004; Wang, Li, Abel, Nakamura 2010; ...)

Stars, especially massive stars, gain most mass by Bondi-Hoyle accretion of ambient clump gas



Originally based on simulations including only thermal pressure.

Massive stars form on the timescale of the star cluster, with relatively low accretion rates.

Massive Star Formation Theories

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t SOFIA Result on Clump Infall Vinfall ~ 0.1 Vff (Wyrowski et al. 2016)

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Violent interactions? Mergers? (Bally & Zinnecker 2005)



Schematic Differences Between Massive Star Formation Theories



The Initial Conditions of Massive Star Formation

Do massive starless cores exist? Are they close to virial equilibrium?





Mid-IR Extinction Mapping of Infrared Dark Clouds

(Butler & Tan 2009, 2012; see also Peretto & Fuller 2009; Ragan et al. 2009; Battersby et al. 2010)

G28.37+00.07

16'



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G28.37+00.07







Formation of IRDCs, GMC Collisions, Dense Gas Mass (Scoville et al. 1986; Tan 2000; Tasker & Tan 2009; Tan 2010; **Fractions & KS Relation** Suwannajak, Tan & Leroy 2014)

Wu, Tan, Nakamura+ (2016)

Importance of [CII] mapping of IRDCs to understand origin of dense gas mass fraction variation in GMCs (Beuther et al.; Ragan et et al.)



Sample of ~50 massive "starless" core/clumps

(Butler & Tan 2012; Butler et al. 2014) Mass surface densities (M=60M_☉)

 $\overline{\Sigma} \simeq 0.1 \text{--} 0.4 \,\text{g}\,\text{cm}^{-2}$

Cores show central concentration $ho \propto r^{-k_{
ho}} \qquad k_{
ho} = 1.5\pm0.3$

Contain many Jeans masses. B-fields suppress fragmentation? Not radiative heating (c.f., Krumholz & McKee 2008).

$$M_{\rm BE} = 1.182 \frac{c_{\rm th}^4}{\left(G^3 P_{s,\rm core}\right)^{1/2}} \to 0.0504 \left(\frac{T}{20 \text{ K}}\right)^2 \frac{1}{\Sigma_{\rm cl}} M_{\odot}$$

Magnetic Critical Mass (Bertoldi & McKee 1992)

$$M_B = 79c_{\Phi}^3 \left(\frac{R}{Z}\right)^2 \frac{\bar{v}_A^3}{(G^3\bar{\rho})^{1/2}} = 1020 \left(\frac{R}{Z}\right)^2 \left(\frac{\bar{B}}{30\,\mu\text{G}}\right)^3 \left(\frac{10^3 \text{ cm}^{-3}}{\bar{n}_{\text{H}}}\right)^2 \quad M_{\odot}$$
$$n_{\text{H}} \sim 10^5 \text{ cm}^{-3}, \text{ B} \sim 200 \mu\text{G} \rightarrow \text{M}_{\text{B}} \sim 100 \text{ M}_{\odot}$$



Four IRDC core/clumps selected to be dark at 8, 24, 70 µm

F2

G2

So use high angular resolution observations of $N_2D^+(3-2)$ to

0.25 0.30 0.35

- 1. Identify exact location of (massive) starless cores
- 2. Measure core velocity dispersion, σ .

- - 1

0.10 0.15

0.20

3. Measure D_{frac}?

C

4. Astrochemical ages?



Comparison to Turbulent Core Model Tan, Kong et al. (2013)



Predictions from Virial Equilibrium

Tan, Kong et al. (2013)

•	1D velocity	dispersion	if virialized:
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$$m_A = \sqrt{3}\sigma_c/v_A = 1)$$

$$\sigma_{\rm c,vir} \rightarrow 1.09 \left(\frac{M_c}{60M_{\odot}}\right)^{1/4} \left(\frac{\Sigma_{\rm cl}}{1\,{\rm g\,cm^{-2}}}\right)^{1/4} \,{\rm km\,s^{-2}}$$

Core	C1-N	C1-S	F1	F2	G2-N	G2-S
Σ_{cl} (g cm ⁻²)	0.48	0.40	0.22	0.32	0.21	0.19
$M_{c}(M_{\odot})$	16	63	6.5	4.7	2.4	0.83
σ _{vir} (km/s)	0.66±0.22	0.88±0.30	0.43±0.15	0.44±0.15	0.33±0.11	0.25±0.09
σ _{obs} (km/s)	0.41±0.03	0.41±0.02	0.25±0.02	0.42±0.04	0.34±0.02	0.30±0.02

 $< \sigma_{obs} / \sigma_{vir} > = 0.81 \pm 0.13$

 $\begin{array}{l} & \longrightarrow & m_{A,vir} = 0.28 \ \text{->} \ B_{vir} = 0.9 \text{mG} \\ & B_{med} \simeq 0.12 n_{H}^{0.65} \ \mu \text{G} \ (\text{for} \ n_{H} > 300 \ \text{cm}^{-3}) \ (\text{Crutcher et al. 2010}) \\ & n_{H,c} = 6.4 \times 10^5 \text{cm}^{-3} \ \text{->} \ B_{med} = 0.7 \text{mG} \end{array}$

Tentative Conclusion: Cores appear to be near virial equilibrium, after accounting for clump envelope. Possibly slightly sub-virial; or have stronger B-fields (see also - Kauffmann, Pillai & Goldsmith 2013).

A Hunt for Massive Starless Cores

Kong, Tan et al. (2016b, arXiv:1609.06008)

- Snapshot ALMA survey of 32 IRDC clumps
- Automated N₂D⁺(3-2) core finding
- ~100 N₂D⁺(3-2) core candidates detected
- Dynamical analysis of 6 best cores: < $\sigma_{obs}/\sigma_{vir}$ > = 0.80 ± 0.06



But are the "Cores" Starless? - sometimes not!

ALMA Cycle 2 follow-up of C1 region



Tan, Kong et al. (2016)

The Deuteration Fraction of C1-S & C1-N

Kong, Tan, Caselli, Fontani, Pillai, Butler, Shimajiri, Nakamura, Sakai (2016) • Multi-transition study of N_2D^+ & N_2H^+



The Deuteration Clock

Kong, Caselli, Tan, Wakelam, Sipilä (2015) (see also Pagani et al. 2009, 2013)

 Modeling of N₂H⁺ deuteration with gas-phase, spin-state network (132 species; 3232 reactions) to constrain age or collapse rate



Parameter Space Exploration: n_H , T, ζ , f_D , OPR^{H2} Deuteration time; comparison with t_{ff} & t_{ad}



The Deuteration Clock

Kong, Caselli, Tan, Wakelam, Sipilä (2015)

Evolving density model



 $\begin{array}{l} \text{If } n_0 \geq 0.1 n_1 \\ \text{If initial OPR}^{H2} \geq 1 \\ \text{initial } f_D = 1 \\ \text{THEN } D_{\text{frac}} \geq 0.1 \Rightarrow \alpha_{\text{ff}} \leq 0.1 \end{array}$



But, observed values of D_{frac} consistent with predicted equilibrium values

 $D_{\text{frac}} \equiv [N_2 D^+]/[N_2 H^+] = 0.15 - 0.72 \text{ (C1-S)} = 0.16 - 0.44 \text{ (C1-N)}$

Magnetized, Turbulent, Massive Starless Core Simulations

Goodson, Kong, Tan et al. (2016, arXiv:1609.07107)

- ATHENA: ideal MHD; isothermal (15K; γ=1.01)
- Parameterized D chemistry (Kong et al. 2015): d[N₂D⁺]/dt (n_H); d[N₂H⁺]/dt (n_H)

C1-S: $M_c=60M_{\odot}$; r=0.07pc; $\Sigma_{cl}=0.5$ g cm⁻²; n_H =6x10⁵ cm⁻³; t_{ff}=40kyr; B₀~2.5mG; B_s~0.5mG





Goodson et al. 2016 - Simulated D_{frac} maps



Chemical Clock with para-H₂D⁺

ortho & para H₂D⁺ to constrain ortho to para ratio of H₂ (Brünken et al. 2014)



Protostellar core IRAS 16293-2422 A/B ($n_H \sim 2x10^5 \text{ cm}^{-3}$, $t_{ff} = 1.0x10^5 \text{ yr}$) OPR_{H2}~10⁻⁴, which indicates chemical processing for >1Myr = 10 t_{ff}

This information helps break degeneracies in Deuteration chemical clocks [N₂D⁺]/[N₂H⁺] (Pagani et al. 2011, 2013; Kong et al. 2015)

Constraints for Initial Conditions of Numerical Simulations

Peters et al. (2011) M = $100M_{\odot}$, R=0.5pc, n_H = 5400cm⁻³, B= 10μ G Seifried et al. (2012) M = $100M_{\odot}$, R=0.25pc, n_H = $4.4x10^4$ cm⁻³, B~1mG

Myers et al. (2013) M = $300M_{\odot}$, R=0.1pc, n_H = $2.4x10^{6}$ cm⁻³, B>~1mG



Do massive protostars have morphologies similar to low-mass protostars? What sets the star formation efficiency from the core? CMF -> IMF?



Protostellar Evolution

Zhang, Tan, Hosokawa (2014)



see also Palla & Stahler 1993; Hosokawa et al. (2010)

Diagnostics of the Turbulent Core Model

Zhang & Tan (2011), Zhang, Tan & McKee (2013), Zhang, Tan & Hosokawa (2014), Tanaka, Tan & Zhang (2016)



Massive Protostar G35.2N: d=2.2kpc; L~10⁵L_o



Spectral energy distribution



Flux profiles along outflow cavity axis



massive star is forming.

10

Temperature, T(K)

The SOFIA Massive (SOMA) Star Formation Survey

Jonathan C. Tan, James M. De Buizer, Mengyao Liu, Yichen Zhang, Jan E. Staff, Maria T. Beltrán, Ralph Shuping, Barbara Whitney

See poster: Mengyao Liu et al.

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THE SOFIA MASSIVE (SOMA) STAR FORMATION SURVEY: I. OVERVIEW AND FIRST RESULTS

VATOR

BSER

ASTRONOM

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Type I: MIR sources in IRDCs - relatively isolated sources in Infrared Dark Clouds, some without detected radio
Type II: Hyper-compact - often jet-like, radio sources, where the MIR emission extends beyond the observed radio emission (e.g., G35.2)
Type III: Ultra-compact - radio sources where the radio emission is more extended than the MIR emission
Type IV: Clustered sources - a MIR source exhibiting radio emission is surrounded by several other MIR sources within ~60"

Also extended to Intermediate-Mass protostars.

The SOMA Survey

20 protostars observed as of Oct 2016 (end of Cycle 4).



Cepheus A



10-6		G35.20-0.7	74		СерА	СерА			NG	C7538 NGC7538	
10 ⁻⁰ 10 ⁻⁷ (- E ⁰ 10 ⁻⁸ - s ⁶	IRAC bands treated as upper limits		clump env subtra 10 ⁻⁷	velope action	SE	D Initial Clum	Core I			sity Sal	
ق 10 ⁻¹⁰ لے 10 ⁻¹¹		0 λ (μm)	• 10 ⁻¹⁰	1000 10	10 100 λ (μm) >	Curre Viewi Foreg	nt Pro ng Ang round	ostella le to O Extinc	utflow tion, A	5, m* Axis , θ _{view} 10 100 (μm) λ (μm)	100 1000
10 ⁻⁶	E	IRAS20126) 		Z	hang &	Tan m	odels			
	Source	χ^2/N_{\sim}	$M_{\mathbf{c}}$	$\Sigma_{ m cl}$	m_*	$\theta_{ m view}$	A_V	M_{env}	$\theta_{w,\mathrm{esc}}$	$\dot{M}_{ m disk}$	$L_{\rm bol}$
- 10 ·			(M_{\odot})	$(g cm^{-2})$	(M_{\odot})	(deg)	(mag)	(M_{\odot})	(deg)	$(M_{\odot}/{ m yr})$	(L_{\odot})
- 5 -10 *	5.20-0.74	.3	120	3.2	12	29	37.6	99	18	9.6×10^{-4}	5.4×10^{4}
ergs		8.0	120	1.0	24	48	57.2	68	37	4.9×10^{-4}	1.5×10^5
لي 10 ⁻¹⁰		8.0	120	1.0	12	29	3.5	96	20	4.0×10^{-4}	5.0×10^4
10		9.8	60	3.2	16	48	81.1	31	32	8.4×10^{-4}	1.2×10^{5}
10 ⁻¹¹	L <u>.</u>	0 10.8	<u>10</u> 60 · · ·		12	48	7.1	38	27	7.6×10^{-4}	5.2×10^4
(Cep A	49 NGC7538	480	0.1	12	83	81.1	458	12	1.0×10^{-4}	2.4×10^4
10 ⁻⁶		5.0	480	0.1	16	89	100.0	441	15	1.2×10^{-4}	3.9×10^{4}
10 ⁻⁷	Ē	6.9	120	0.3	12	62	61.4	93	24	1.6×10^{-4}	3.7×10^4
cu3		7.0	60	3.2	16	68	87.0	31	32	8.4×10^{-4}	1.2×10^{5}
		7.4	120	1.0	24	55	100.0	68	37	4.9×10^{-4}	1.5×10^{5}
Ů, No C	C 7538	0.6	480	0.1	16	22	9.3	441	15	1.2×10^{-4}	3.9×10^4
م 10 ⁻¹⁰		1.2	240	0.1	24	44	37.6	171	33	1.1×10^{-4}	8.3×10^4
10-11		1.4	240	Q.1	32	48	65.8	140	42	1.1×10^{-4}	1.5×10^{5}
10	1 1	$0 1.7_{\lambda \ (\mu m)}$	160	13002	12	34	14.2	38	27	7.6×10^{-4}	5.2×10^{4}
		2.3	60	3.2	16	39	61.4	31	32	8.4×10^{-4}	1.2×10^{5}

Feedback During Massive Star Formation

Is there a maximum stellar mass set by by formation processes?



m∗ Feedback processes:

- 1. Protostellar outflows
- 2. Ionization
- 3. Stellar winds
- 4. Radiation pressure
- 5. Supernovae





m_{*max}~150 M_☉ (e.g. Figer 2005).

But Crowther et al. (2010) claim most massive star to form was initially ~300M_☉, consistent with statistical sampling of Salpeter IMF with no maximum cutoff mass.

Staff+ (2010); Kuiper+ (2015)

Peters et al. 2010, 2011

Krumholz+ (2009); Rosen+ (2016) Kuiper et al. (2012); Klassen+ (2016)

Accretion processes: Core/disk fragmentation (Kratter & Matzner 06; Peters et al. 10) Stellar processes: Nuclear burning instabilities/enhanced mass loss Currently unclear what sets the shape of the massive star IMF

Dynamical Interactions: Massive Protostars in Crowded Environments



Orion KL protostar perturbed by a passing runaway star (BN) ejected from the Trapezium star $\theta^{1}C$

Chatterjee & Tan (2012)



Orion KL: a perturbed massive star-forming core

Tan (2004) Chatterjee & Tan (2012)

HOWEVER, SEE: Bally & Zinnecker 2005; Rodriguez et al. (2005) Gomez et al. (2005) **Gomez et al. (2008)** Zapata et al. (2009) Zapata et al. (2011a,b) **Bally et al. (2011)** Goddi et al. (2011) Moeckel & Goddi (2011)

BUT. SEE Plambeck et al. (2016)

interplay of stellar and gas dynamics. CISCO (H2 (v=1-0 S(1)) - Cont) stronomical Observatory of Japan January 28, 1999

Dynamical "harassment" of

protostellar cores in star clusters.

Star cluster formation involves

θ¹C

Massive Star Formation Theories:ConclusionsCore Accretion; Competitive Accretion; Protostellar Collisions

Theory: "Turbulent Core Model": normalize core surface pressure to surrounding clump pressure, i.e. self-gravitating weight. Core supported by nonthermal pressure (B-fields/turbulence).

1: Massive starless/early-stage cores exist

2: They are near virial equilibrium

3: Massive protostars can have a similar morphology to low-mass protostars, but dynamical interactions can occur (BN/KL)

4. SOFIA is playing a crucial role: IRDC formation, clump infall, astrochemical ages, 0 SEDs and images, and more!





